

**The Fujita Exponent**  
for  
**Semilinear Heat Equations**  
with  
**Quadratically Decaying Potentials**  
or  
**in an Exterior Domain**

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Consider the classical semilinear heat equation

$$\begin{aligned}u_t &= \Delta u + u^p \text{ in } R^n \times (0, T); \\u(x, 0) &= \phi(x) \geq 0 \text{ in } R^n,\end{aligned}\tag{1}$$

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- ▶ The above result goes back to **Fujita ('66)** in the case  $p \neq p^*$ . Various proofs of blow-up in the borderline case  $p = p^*$  can be found in *Aronson and Weinberg ('76)*, *Kobayashi, Siro and Tanaka ('77)* and *RP ('97)*.

More recently, **Zhang ('01)** considered the consequences of adding a potential term:

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► **Theorem (Zhang)** *Let  $n \geq 3$ .*

*i. If  $0 \leq V(x) \leq \frac{\omega}{1+|x|^b}$ , with  $b > 2$  and  $\omega > 0$ , then  $p^* = 1 + \frac{2}{n}$ , and consequently the potential does not affect the critical exponent;*

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- ▶ **He also noted that it is unclear whether or not  $p^*$  is finite in the case that  $V(x) \sim \frac{\omega}{|x|^2}$ , with  $\omega < 0$ .**

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- ▶ Since we are assuming here that  $\omega > 0$ , one has  $\alpha(\omega, n) > 0$  and thus  $p^*(\omega) < 1 + \frac{2}{n}$ .

$$\alpha(\omega, n) = \frac{2 - n + \sqrt{(n-2)^2 + 4\omega}}{2}, \quad \omega > 0$$

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- If  $V(x) \geq \frac{\omega}{|x|^2}$ , for large  $|x|$ , then for  $p > p^*(\omega)$  there exist global solutions;*
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- ▶ Ishige's proof involved comparison with a solution to the radially symmetric linear equation  $v_t = \Delta v - \hat{V}(|x|)v$ , where  $\hat{V}(r) \sim \frac{\omega}{r^2}$  as  $r \rightarrow \infty$ . The **asymptotic space-time behavior** of this equation appears in *Ishige and Kabeya (08)*.

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- ▶ **What happens when  $\omega < 0$ ?**

Our method will require us to consider more generally

$$\begin{aligned}u_t &= \Delta u - Vu + au^p \text{ in } R^n \times (0, T); \\u(x, 0) &= \phi(x) \geq 0, \text{ in } R^n,\end{aligned}\tag{2}$$

where

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- ▶ But now we want to consider (2) with  $V(x) \sim \frac{\omega}{|x|^2}$  for large  $|x|$  with  $\omega < 0$  and also with  $\omega > 0$ .

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**Ishige's result: when  $a \equiv 1$  and  $V(x) \sim \frac{\omega}{|x|^2}$  with  $\omega > 0$ , one has  $p^* = 1 + \frac{2}{n+\alpha(\omega, n)}.$**

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### ► Theorem

Let  $-\frac{1}{4}(n-2)^2 \leq \omega < 0$ ,  $n \geq 3$ .

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► **Remark.** We also have a similar result for  $\omega > 0$ , with  $n \geq 2$ .

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▶ **Theorem (Zhang)** iii. If  $\frac{\omega}{1+|x|^b} \leq V(x) \leq 0$ , for  $b > 2$  and  $\omega < 0$  with  $|\omega|$  sufficiently small, then  $p^* = 1 + \frac{2}{n}$ ;

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$$\begin{aligned} u_t &= \Delta u - Vu + au^p \text{ in } R^n \times (0, T); \\ u(x, 0) &= \phi(x) \geq 0, \text{ in } R^n. \end{aligned} \tag{6}$$

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*If there exists a bounded domain  $D \subset R^n$  for which the principal eigenvalue  $\lambda_{0;D}(-\Delta + V) < 0$ , then there are no global solutions to (6) for any  $p > 1$ ; that is,  $p^* = \infty$ .*

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*If  $\inf \text{spec}(-\Delta + V; R^n) < 0$ , or equivalently, if there exists a bounded domain  $D \subset R^n$  for which the principal eigenvalue  $\lambda_{0;D}(-\Delta + V) < 0$ , then there are no global solutions for any  $p > 1$ ; that is,  $p^* = \infty$ .*

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Let  $n \geq 2$ . Let  $V(x) \leq \frac{\omega}{|x|^2}$  with  $\omega < -\frac{(n-2)^2}{4}$ , for  $|x| > \epsilon$  and  $\epsilon > 0$  sufficiently small. Then there are no global solutions to (7) for any  $p > 1$ ; that is,  $p^* = \infty$ .

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$$u_t = u_{rr} + \frac{n-1}{r} u_r - V(r)u + a(r)u^p, \quad (8)$$

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- ▶ So the **solution to (12) is dominated by the solution to**

$$w_t = w_{rr} + \frac{N-1}{r} w_r + \hat{a}(r) w^p \text{ in } (r_0, \infty) \times (0, T);$$

$$w(r_0, t) = 0, \quad t > 0,$$

where

$$N \equiv n + 2\alpha(\omega, n) \geq 2,$$

$$\hat{a}(r) = r^{\alpha(p-1)} a(r),$$

$$c_1 r^M \leq \hat{a}(r) \leq c_2 r^M, \text{ for large } r,$$

where  $M \equiv \alpha(\omega, n)(p-1) + m$ .



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► Recall that for

$$u_t = \Delta u + au^p \text{ in } R^n \times (0, T);$$

$$c_1 |x|^m \leq a(x) \leq c_2 |x|^m, \text{ for large } |x|,$$

one has  $p^* = 1 + \frac{(2+m)^+}{n}$ .



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- ▶ **So it is reasonable to suspect that no global solution exists for (13) if**

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- ▶ Of course,  $N$  is in general fractional and also, we have placed the Dirichlet b.c. at  $r = r_0$  which serves to make solutions smaller.



$$w_t = w_{rr} + \frac{N-1}{r} w_r + \hat{a}(r) w^p \text{ in } (r_0, \infty) \times (0, T); \quad (14)$$

$$w(r_0, t) = 0, \quad t > 0,$$

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**It is reasonable to suspect that no global solution exists for (14) if**

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**It is reasonable to suspect that no global solution exists for (14) if**

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▶ But (15) is equivalent to

$$p \leq 1 + \frac{(2+m)^+}{n + \alpha(\omega, n)}.$$



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**It is reasonable to suspect that no global solution exists for (14) if**

$$p \leq 1 + \frac{(2+M)^+}{N}. \quad (15)$$

▶ But (15) is equivalent to

$$p \leq 1 + \frac{(2+m)^+}{n + \alpha(\omega, n)}.$$

▶ Or equivalently,  $p \leq p^*(\omega, m)$ .



$$\begin{aligned} u_t &= \Delta u + au^p \text{ in } R^n \times (0, T); \\ c_1|x|^m &\leq a(x) \leq c_2|x|^m, \text{ for large } |x|; \end{aligned} \tag{16}$$

There are no global solutions for  $p \leq 1 + \frac{(2+m)^+}{n}$ .

$$w_t = w_{rr} + \frac{N-1}{r}w_r + \hat{a}(r)w^p \text{ in } (r_0, \infty) \times (0, T);$$

$$w(r_0, t) = 0, \quad t > 0;$$

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**Need to show that there are no global solutions for**

$$p \leq 1 + \frac{(2+M)^+}{N}. \tag{17}$$



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**Need to show that there are no global solutions for**

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- ▶ The proof of (16) uses heavily the explicit form of the **Gaussian kernel** for the linear part of (16). To use that type of proof for (17) requires use of heat kernel  $q_{N,r_0}(t, r, \rho)$  for linear part of (17).

▶

$$u_t = \Delta u + au^p \text{ in } R^n \times (0, T);$$

$$\text{No global solutions for } p \leq 1 + \frac{(2+m)^+}{n}. \quad (18)$$

$$w_t = w_{rr} + \frac{N-1}{r}w_r + \hat{a}(r)w^p \text{ in } (r_0, \infty) \times (0, T);$$

$$w(r_0, t) = 0, \quad t > 0;$$

**Need to show no global solutions for**  $p \leq 1 + \frac{(2+M)^+}{N}$ .

$$(19)$$

Proof of (18) uses the explicit form of the **Gaussian kernel** for the linear part of (18).

▶ 
$$u_t = \Delta u + au^p \text{ in } R^n \times (0, T);$$

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Proof of (18) uses the explicit form of the **Gaussian kernel** for the linear part of (18).

▶ Let  $q_{N,r_0}(t, r, \rho)$  be kernel for linear part of (19).

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$$\text{No global solutions for } p \leq 1 + \frac{(2+m)^+}{n}. \quad (18)$$

$$w_t = w_{rr} + \frac{N-1}{r} w_r + \hat{a}(r)w^p \text{ in } (r_0, \infty) \times (0, T);$$

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$$(19)$$

Proof of (18) uses the explicit form of the

**Gaussian kernel** for the linear part of (18).

- ▶ Let  $q_{N,r_0}(t, r, \rho)$  be kernel for linear part of (19).
- ▶ The corresponding heat kernel on **whole space** is

$$q_N(t, r, \rho) \equiv \exp\left(-\frac{r^2 + \rho^2}{4t}\right) \frac{\rho^{N-1}}{2t(r\rho)^{\frac{N}{2}-1}} I_{\frac{N}{2}-1}\left(\frac{r\rho}{2t}\right),$$

where  $I_\nu$  is the modified Bessel func. of order  $\nu$ .

▶

$$u_t = \Delta u + au^p \text{ in } R^n \times (0, T);$$

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**Need to show no global solutions for**  $p \leq 1 + \frac{(2+M)^+}{N}$ .

$$(21)$$

Let  $q_{N,r_0}(t, r, \rho)$  be kernel for linear part of (21).

The corresponding heat kernel on **whole space** is

$$q_N(t, r, \rho) \equiv \exp\left(-\frac{r^2+\rho^2}{4t}\right) \frac{\rho^{N-1}}{2t(r\rho)^{\frac{N}{2}-1}} I_{\frac{N}{2}-1}\left(\frac{r\rho}{2t}\right).$$

**If we can use  $q_N(t, r, \rho)$  instead of  $q_{N,r_0}(t, r, \rho)$ , then we can prove (21) via method of (20).**

$$u_t = \Delta u + au^p \text{ in } R^n \times (0, T);$$

$$\text{No global solutions for } p \leq 1 + \frac{(2+m)^+}{n}. \quad (20)$$

$$w_t = w_{rr} + \frac{N-1}{r} w_r + \hat{a}(r)w^p \text{ in } (r_0, \infty) \times (0, T);$$

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► **Need:**

$$q_{N,r_0}(t, r, \rho) \geq cq_N(Kt, r, \rho), \quad r, \rho \geq r_0 + 1, \quad \text{for some } K > 0.$$

▶

$$w_t = w_{rr} + \frac{N-1}{r} w_r \text{ in } (r_0, \infty) \times (0, T); \quad (22)$$

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- ▶ Recall:  $\alpha(\omega, n) = \frac{2-n+\sqrt{(n-2)^2+4\omega}}{2}$ ,  $\omega \geq -\frac{(n-2)^2}{4}$ .

**So  $N > 2$  unless  $\omega = -\frac{(n-2)^2}{4}$ .**

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- ▶ For  $N > 2$ , the Markov diffusion process corresponding to the generator  $\frac{d^2}{dr^2} + \frac{N-1}{r} \frac{d}{dr}$  is **transient**. Using this and the **parabolic Harnack inequality**, **Grigor'yan and Saloff-Coste ('02)** proved that

$$q_{N,r_0}(t, r, \rho) \geq cq_N(Kt, r, \rho), \text{ for } r, \rho > r_0+1, \text{ for some } K > 0.$$

► **Fujita Exponent in an Exterior Domain**

Let  $B_1 = \{x \in R^n : |x| < 1\}$ ,

$$\begin{aligned} u_t &= \Delta u + au^p \text{ in } (R^n - \bar{B}_1) \times (0, T); \\ u(x, t) &= 0, \quad |x| = 1, \quad t \geq 0; \\ u(x, 0) &= \phi(x) \geq 0, \text{ in } R^n - \bar{B}_1, \\ \text{where } c_1|x|^m &\leq a(x) \leq c_2|x|^m, \text{ for large } |x|. \end{aligned} \tag{23}$$

$$\text{Let } p^* = 1 + \frac{(2+m)^+}{n}$$

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► **Theorem**

Let  $n \geq 2$ .

- i. If  $1 \leq p \leq p^*$ , then there exist global solutions to (23);*
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## ► Fujita Exponent in an Exterior Domain

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► For proof of the theorem, we use estimates on heat kernel for

$u_t = \Delta u$  in an exterior domain, obtained by **Grigor'yan and Saloff-Coste ('02).**